# A study on Vortex Induced Vibrations Suppression Devices for Marine Pipeline 

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## ABSTRACT

Vortex induced vibration is caused due to the periodic shedding of vortices from either sides due to flow separation over a bluff body ${ }_{\lrcorner}^{1}$ An asymmetrical flow pattern develops around the body due to the periodic shedding of vortices and changes the pressure distribution on the surface of the body which creates fluctuating lift forces on either side of the body, thus making the structure to vibrate. The frequency of shedding depends on diameter-of the body and the velocity of flow (Reynolds Number). High amplitudes of vibrations can occur when the frequency of vortex shedding coincides with the natural frequency of the structure Yortex induced vibrations ( V F ) are an important source for fatiguedamage of the structures. This phenomenoncan typically occur on long chimneys, cables used in bridges, power transmission lines in air. In marine environment the structures subjected to these vibrations are marine risers, anchor lines and free spanning pipelines etc, Large diameter-pipes used todraw cold water from deeper water depths could be susceptible to vortex induced vibrations. CurrentlyNIOT is working on a bundle of pipes forming a conduit. Each pipe is made of HighdensitityPolyethylene (HDPE). The bundles with 2 and 3 cylinders are considered for the study. The flow around these pipes will be studied by performing CFD analysis considering the cylinders to be stationary. The present work is focused on studying flow behaviour around the pipes, vortex sheddingaround cylinders using-Ansys Fluent:- -

KEYWORDS: Vortex Induced Vibration (VIV); Computational Fluid Dynamics (CFD); Ansys Fluent:

## CHAPTER 1

## INTRODUCTION

### 1.1 General

Vortex-shedding is phenomen that occurs over bluff bodies subjected to sustained currents, which might result in large transverse motions (Humphries and Walker 1987). When a flowing medium strikes a nonstreamlined bluff object, it moves around the object generating alternating pressure forces on either side of the bluff body. This results in the formation of vortices, causing periodic forces, which are strong enough to set the body into oscillatory motion. Vortex-Induced Vibration (VIV) are the outcome of such exciting forces that are generated by vortex-shedding on hull of a bluff body, causing response closer to the resonant period (Anagnostopoulos and Bearman 1992; Bearman 1984). Experimental investigations carried out on elastically mounted cylinder showed wake formation with the maximum amplitude occurring near lower limit of lock-in region. Due to VIV, structure undergoes a number of stress cycles leading to fatigue damage. Based upon the
significant consequences of VIV on several mechanical systems and ocean engineering structures, VIV suppression configurations are encouraged (Sarpkaya 1978). Observation of large amplitude responses of systems that used water as working fluid necessitates the importance of VIV suppression by design (Khalak and Williamson 1991; Zdravkovich 1981). While several suppression techniques of VIV suppression are successfully developed, cost and difficulty of implementation restricted their applications (Owen and Masa 2003). In view of suppressing VIV, use of helical strakes proved to be effective in strengthening the cylinder to resist larger bending moments that result from the increased drag. Under close examination of flow behind a pair of cylinders, pair of anti-phase streets and in-phase shedding can cause the formation of a single wake in large scale (Govardhan and Williamson 2000; Brika and Lanevilfe 1993). Their applications in marine risers by modelling them, as flexible cylinders are quite successful. 3987 Applications in Preliminary Analysis and Design Normalized vortex'shedding frequency, which is represented by Stroublal fumber is affected by various flow velocities (Lesage and-Garthshore 1987). An application example on experimentalinvestigations of a rigid cylinder, resembling hall of a Spar platform is presented. Flow around a cylinder is one of the classical topics in field of hydrodynamics (Mutlu and Jorgen 2003). Models proposed by several researchers consider rigid circular cylinder with a single degree of freedom in the cross-flow direction (Gabbai and Benaroya 2005; Rodolfo et al. 2011). Experimental investigations carried out on long, flexible cylinders indicated the formation of vortex-induced motion in the form of hysteresis loop motion in which contributions are seen from of the vortex shedding modes (Chandrasekaran and Marin 2016). Amongst various methods proposed to-suppress VIV but showed a substantial increase in drag, resulting in high drift. When strakes or bumps are fixed, no regular shedding is observed on the wake side of the cylinder, while major advantage is effective suppression of amplitude of vibration of the cylinder. Recent studies also show the possibilities of suppression of vibration of mutiotecylinders using rough strip (Blevins 1994). A few analytical and numerical investigations of VIV on cylinders with low mass and aspect ratio are helpful in modelling wake oscillations and force decomposition, resulting from VIV suppression systems. Alternatively, force reduction is also achieved by providing an outer perforated cover, which alters fluid flow around cylinders in flow regime.

### 1.2 Objective of the Study

- To study on the Jortex-Suppression Devices for marine Pipetines.


### 1.3 Scope of the Stuidy

- To study the flow behavior around a bare cylinder for different current speeds using CFD.
- To study the flow behavior around a cylinder with different vortex suppression devices (VSD) for different current speeds using CFD.
- Analysis of flow around a cylinder with optimized VSD for different current speeds.


### 1.4 Literature Review

1. Meghanadhan C A and Sunil A S (2016) worked on study of Vortex Shedding Behind a Bluff Body in a Water Channel. They conducted experiment in water channel installed in fluid mechanics lab to study the flow characteristics and flow pattern in the wake region of a circular cylinder and understand the relation between the Reynolds number and Strouhal number. They found that there was no vortex shedding at $\mathrm{Re}=40$. And by the experiment they concluded that the cylinder with helical stakes was the effective device for suppression of vortex induced vibrations as passive control. And also, frequency of vortex shedding was reduced by helical stakes of suitable pitch and height.

2. Alok Mishra and Ashoke de studied-on suppression of vortex shedding using slit through a circular cylinder at a low Reynolds number. They performed 2-D simulations on bare cylinder and cylinder with a slit along with the incoming flow. The investigation was done for different slit ratios ( $\mathrm{S} / \mathrm{D}=0$ to 0.25 ) at different Re values ( $\operatorname{Re}=100$ to $50 \theta$ ) to find the slit angle $(\theta)$ on vortex shedding. The $\mathrm{S} / \mathrm{D}=0.25$ at $\operatorname{Re}=500$ with slit angle $\left(\theta=20^{\circ}\right)$ illustrates the maximum suppression in vortex shedding.
3. Zheng-Shou Chen and Wu-Joan Kim (2010) have done a Numerical investigation of vortex shedding and vortex-induced vibration for flexible riser models. They used ANSYS MFX solver to solve the twoy wh problems related to VIV of flexible riser system. They carried out simulations for two flexible riser model systems, one was related to the simulation of vortex visualization in the wake for a riser model subject to forced oscillation, and another was related to the simulation of fluid-structure interaction between the pipes of coupred multi-assembled riser system. The riser model was made of PVC pipes with an outer diameter of 50 mm and wall thickness of 2.5 mm . By the study they concluded that in the multi-assembled pipe systems, the external current velocity and the distance between two pipes are the critical factors to determine the vibration state and the steady yibration state in quad-pipe system may be destroyed more easily than dual-pipe system.

4. Dwi Priyanta and Anang Abdullah carried out Analysis of the Suppression Device as Vortex Induced Vibration (VIV) Reducer on Free Span using Finite Element Method. In this study they have used twe types of Suppression deyife i.e., Helical strike and fairing to compared in reduce the impact of YIV They have done fatigue life calcuation for the happening VIV process simulation of Finite Element Method (FEM) using software Ansys Flüent. They concluded that assumed operational designed time of the pipe-for 50 years, by the test they found that on the condition of the critical free span exposed VIV the life time value on plain pipe 44.21 years, at and Fairing pipe with 53.09 years and on the pipe with the Helical 52.95 strike year.
5. T. Zhou, S.F. Mohd. Razali, Z. Hao, L. Chenga studied on the study of vortex-induced vibration of a cylinder with helical strakes. They used 80 mm cylinder and helical strakes of dimensions of 10 d in pitch and 0.12 d in height. It was tested in a wind tunnel and they found that the helical strakes can reduce VIV by about $98 \%$.
6. Satya Prakash Singh, Goutam Biswas, Perumal Nithiarasu have done a numerical study of vortex shedding from a circular cylinder vibrating in the in-line direction. In this study a circular cylinder, vibrating in the inline direction is placed in a fluid stream and assumed forced vibration as sinusoidal and considered distance between the mean centre of the cylinder and upstream and downstream boundaries are, respectively, 10Dand 67.5 D , where D is the diameter of the cylinder. They observed that the drag and lift coefficients were perfectly periodic at a Reynolds number of 200 and they reach a chaotic pattern as the Reynolds number is increased to a value of 350 .
7. Nur Ain Shafiza Ramzi, Lee Kee Quen, Hidetaka Senga, Hógi-Siáng Kang, Meng Hee Lim, Noor Idora Mohd Sukarnoor these all people studied on Suppression of vortex indưced, vibration of a rigid cylinder using flexible shrouding. They conducted experimental investigations to find the reduction of VIV and mean drag coefficient (Cd mean) acting on the oscillating cylinder fitted with a shroud of various parameters. By the study they observēd that the shroud with mesh size $B=0.2 \mathrm{D}$ and mesh thickness $\mathrm{H}=0.03 \mathrm{D}$ achieyed an optimum condition where it was able to suppress the amplitude without increasing Cd mean, with $76.62 \%$ of amplitude reduction and $48.40 \%$ of Cd mean reduction.

## CHAPTER 2

## THEORETICAL BACKGROUND

### 2.1 General <br> 

In this chapter, the theoretical fluid behaviour around cylinder is included. The details of software tool ANSYS FLUENT are also included. A brief description about the graphical user interface of the software tool is given. The details of various turbulence models available with ANSYS FLUENT software are included.

### 2.2 VIV Theory

When arstructure, in this case, a spar platform is subjected to the fluid flow, somehow the mightexperience excitations or vibrations. These vibrations known as the flow induced vibrations can lead to the fatigued damage to the structure. Hence, it is essential to take these vibrations into consideration whilst designing structures, particularly cylindricalestructures.

### 2.2.1 Regime of flow

One of the non-dimensionless hydrodynamic numbers that is used to describe the flow around a smooth cylinder is the Reynolds number (Re). By the definition, the Reynolds number is the ratio of the inertia forces to viscous forces and formulated as

$$
\begin{equation*}
R e=\frac{U D}{\gamma} \tag{2.1}
\end{equation*}
$$

Where D is the diameter of cylinder, U is the flow velocity and $\gamma$ is the kinematic viscosity of the fluid. Flow regimes are obtained as the result of tremendous changes of the Reynolds number. The changes of the Reynolds number create separation flows in the wake region of the cylinder, which are called vortices. At low values of $\operatorname{Re}(\operatorname{Re}<5)$, there no separation occurs. When the Re is further increased, the separation starts to occur and becomes unstable and initiates the phenomenon called vortex shedding at certain frequency. As the result, the wake has an appearance of a vortex street as can be seen in Table 2.1


Figure 2.1 Flow regime around smooth circular cylinder in steady current. (sumer, 2006)

### 2.2.2 Vortex Shedding

The vortex shedding phenomenon appears when pairs of stable vortices are exposed to small disturbances and become unstable at Re greater than 40 . For these values of Re, the boundary layer over the cylinder surface will separate due to the adverse pressure gradient imposed by the divergent geometry of the flow environment at the rear side of the cylinder. As mentioned in the previous section, vortex shedding occurs at a certain frequency, which is called as vortex shedding frequency (fv). This frequency normalized with the flow velocity $U$ and the
cylinder diameter D, can basically be seen as a function of the Reynolds number. Furthermore, the normalized vortex-shedding frequency is called Strouhal number (St), and formulated as:

$$
\begin{equation*}
S t=\frac{f v D}{U} \tag{2.2}
\end{equation*}
$$

The relationship between Re and St can be shown in Figure 2.2


The large increase in St at the supercritical region is caused by the delay of the boundary separation. It is known that the separation point of the subcritical regime is different from that of the supercriticalregime as shown in Figure 2.2. At the supercritical flow regime, the boundary layers on both sides of the cylinder are turbulent at the separation point. Consequently, the boundary layer separation is delayed since the separation point moves downstream. At this point, the vortices are close to each otherand create faster rate than the rate in the subcritical regime, thereby leading to higher values of the Strouhal number.


Figure 2.3 Separation point of the subcritical regime and supercritical regime.

When Re reaches the value of $1.5 \times 10^{6}$, the boundary layer completely becomes turbulent at one sideand laminar at the other side. This asymmetric situation is called the lee-wake vortices. What happensnext is that lee-wake vortices inhibit the interaction of these vortices, resulting in an irregular and disorderly vortex shedding. When Re is increased to values Turbulent separation in supercritical regime Turbulent separation in subcritical regime 13 larger than $4.5 \times 10^{6}$ (transcritical regime), the regular vortex shedding is re-established and St takes the values of $0.25-0.30$.

### 2.2.3 Drag and Lift Forces

As the result of the periodic change of the, vortex shedding, the pressure distribution of the cylinder due to the flow will also change periodically thereby generating a periedic variation in the force components on the cylinder. The force components can be divided into crossflow and in-line directions. The force of the crossflow direction is commonty named as the lift force (FL) while the latter is named as the drag. force (FD). The lift force appears when the vortex shedding starts to occurand it fluctuates at the vortex shectling frequency. Similarly, the dragforce also has the oscillating partdue to the vortex shedding, but in addition it also hals a force as a resultof friction and pressure difference; this part is called the mean drag. Both of the lift and drag forces are formulated as follows:

$$
\begin{align*}
& F_{L}=\hat{F}_{L} \sin \left(\omega_{s} t+\varphi_{s}\right)  \tag{2.3}\\
& F_{D}=F_{D} \sin \left(2 \omega_{s} t+\varphi_{s}\right) \tag{2.4}
\end{align*}
$$

FL and F D are the amplitudes of the oscillating lift and drag respectively and is the mean drag. Itshould be noted that drag force oscillated with twice the frequency of oscillation of lift force.

### 2.2.4 Vortex Induced Motions

Wheneven any offshore component like spar hull, pipe, risers are open to a steady current, vortices andalternating eddieswill be formed. A vortex forms in the re-circulation to equalise pressure. Disturbances in the flow causes the vortex to be shed, creating a new vortex to equalise the pressure from the preceding vortex. These vertices shed at a regular interval of period. This phenomenon is referred to as vortex shedding. Vortex shedding is an effect of Viscosityland Reynolds number. Reynolds number is a dimensionless quantity which relates which relates inertial forces to the viscousforces in the flow and is given by

$$
\begin{equation*}
\bar{R}_{e}=\frac{\rho u D}{\mu} \tag{2.5}
\end{equation*}
$$



Where, $\rho$ is the density, u is the velocity, D is the diameter and $\mu$ is the dynamic viscosity.
At low Reynolds number, viscous forces dominate the flow while at high Reynolds number, inertial forces dominate. The vortex shedding frequency fs, which is also called as Strouhal frequency is related to a nondimensional Strouhal number, S

$$
\begin{equation*}
f_{s}=\frac{S V_{C}}{D} \tag{2.6}
\end{equation*}
$$

Where, Vc is current velocity and D diameter of cylinder.

Strouhal shedding period is defined as the inverse of the Strouhal shedding frequency.

$$
\begin{equation*}
T_{s}=\frac{1}{f_{s}} \tag{2.7}
\end{equation*}
$$

Natural frequency of structure fn,


Lock-in phenomena is the tendency of the natural period of the structure to match or to synchronize with the shedding period of the structure, where the cylinder will try to lock-in to the frequency shed by the vortices.


$$
\begin{equation*}
s=0.198\left(1-\frac{19.7}{\operatorname{Re} D}\right) \tag{2.10}
\end{equation*}
$$

The Strouhal number of a static pipe or a cylinder is a function of Reynolds number, but it is less dependent on the surface roughness of the body and the turbulence of the free stream. In the Reynoldsnumber range $250<\mathrm{Re}$ $<2 \times 1 \theta^{5}$ the empirical formula can be used.


### 2.2.5 Reduced velocity

When a structure' starts to vibrate due to VIV, the reduced velocity becomes useful. The reduced velocity gives information whin velocity ranges VIV will be maximum. It has been seen that the vibration amplitude reaches maximum for redaced velocities which are in the "lock-in" range. For subcritical flow leckin occurs at reduced velocities $\mathrm{Ur}=5$ to 7 approximately. Reduced-velocity canbe divided into two types, true reduced velocity Ur true and nominal reduced velocity Ur nom. The true reduced velocity is based on the frequency at which the cylinder is actually vibrating whilst the nominal reduced velocity is based on the nominal natural frequency, e.g., natural frequency in still water. Reduced velocity is formulated as follows.

$$
\begin{equation*}
U_{r}=\frac{U}{f_{n} D} \tag{2.11}
\end{equation*}
$$

## CHAPTER 3

## CFDAND NUMERICALSTUDY

### 3.1 General

The details of software tool ANSYS FLUENT are also included. A brief description about the graphical user interface of the software tool is given. The details of various turbulence models availablewith ANSYS FLUENT software are included. Methodology used such $\overrightarrow{\text { as }}$ gémétric mödelling, meshing, physics modelling, solver set up and stopping criteria is explained here.

### 3.2 Computational Fluid Dynamics

Computational fluid dynamics, usually abbreviated as CFD, is a computer-based simulation ornumerical 27odelling of fluid mechanics to solve and 27odelli problems related to fluid flows, heat transfer and associated phenomena such as chemical reactions. CFD provides wide range application for many industrial and nonindustrial areas CFD is a very powerful technique regarding thesimulation of fluid flows.

The application of CFD began in 1960's when the aerospace industry has integrated CFD technique into the design, R \& D and manufacturing of aircrafts and jet engines. CFD codes are being accepted as design tools by many industrial users. Today, many industrials, for instance, ship industry, power plant, machinery, efectronic engineerihg, chemical process, marine engineering and environmental engineering use CFD as a one of the best designtools. Unlike the model testing facility or experimentallaboratory, in CFD simulations there is no need fora big facility. Furthermore, CFD also offers no capacity limit, no model scale limit and cost and costtand schedule efficiency. Indeed, the advantagesby using CFD compared to experimental-based approaches can be concluded as follows:
?

a. Ability to assess a system that controls experiments is difficult or impossible to perform (very largesystem).
b. Ability to assess a system under hazardous conditions (e.g., safety study and accident investigation):
C. Gives unlimited level of results.

This article describes basic theories and information of fluid dynamics and CFD. Given the ries and information are delivered in brief and general. More details of these can be found out at related sources.FEUIDTHEORY

As mentioned in previous section, CFD is the science of predicting fluid flow, heat and mass transfer,chemical reactions, and other related phenomena. The CFD problems are stated in a set of mathematicalequations and are solved numerically. These set of mathematical equations are based on the conservation laws of fluid motion, which are conservation of mass, conservation of momentum, conservation of energy etc. For CFD related to fluid flow, the set of mathematical equations are based on conservation of mass and momentum.

## Conservation of mass

The mass conservation theory states that the mass will remain constant over time in a closed system. This means that the quantity of mass will not change and, the quantity is conserved.
The mass conservation equation, also called the continuity equation can be written as:

$$
\begin{equation*}
\frac{\partial \rho}{\partial t}+\frac{\partial(\rho u)}{\partial x}+\frac{\partial(\rho v)}{\partial y}+\frac{\partial(\rho w)}{\partial z}=S_{m} \tag{3.1}
\end{equation*}
$$

Equation (3.1) is the general form of the mass-conservation equation and is valid for incompressibleas well as compressible flows. The sodrge Sm is the mass added to the system and any user-definedsources. The density of the fluid is $\rho$ and the flow of mass in $x, y$ and $z$ direction is $u, v$ and $w$.

## Conservation of momentum.



The conservation of momentum is originally expressed in Newton's second law. Like the velocity, monentum is a vector quantity as well as a magnitude. Momentum is also a conserved quality, meaning that for a coosed system, the total momentum will not change if there is no external force.
The momentum conservation equation also known as Navier-Stokes equation in general vector form can be written as:


$$
\begin{equation*}
\rho\left(\frac{\partial \bar{v}}{\partial t}+(\bar{v} . \nabla) \bar{v}\right)=-\nabla p+\nabla^{2} \bar{v}+\rho f \tag{3.2}
\end{equation*}
$$

### 3.3.ANSYS FLUENT

ANSYS FLUENT delivers a powerful CFD tool in an easy-to-use environment. ANSYS FLUENT is based on object-oriented programming technology. It is designed to handle large models quickly and efficiently asing a unique client server architecture that seamlessly meshes and simultaneously solves and post-processes over multiple eormputing resources without requiring additional effort fromthe user. Much more than just a CFD solver, ANSYS FLUENT is an entire engineering process for solving problems involving ftow (of fluids or solids), heat transfer, and stress. The software uses a finite volume approach to solve the problems involving flow of fluids.


### 3.3.1 GENERAL WORKFLOW OF ANSYS FLUENT



Túrbulent flow can be defined as a chaotic, fluctuating and randomly condition of flow, i.e. velocity fields. These fluctuations mix transported quantities such as momentum, energy, and species concentration, and eause the transported quantities to fluctuate as well. Turbulence is a time-dependentprocess. Three basic approaches to modelling turbulence are available in Ansys Fluent. These are DNS(direct numerical solution), RANS (Reyrolds averaged Natier-Stokes), and LES (large eddy simulation).


## Reynolds Averaged Navier-Stokes (RANS)

RANS equations are the time-averaged equations of motion of fluid flow. They govern the fransport of the averaged flow quantities, with the complete range of the turbulent scales being modelled. Therefore, it greatly reduces the required computational effort and resources and is widely adopted for practical engineering applications. To obtain the Reynolds-Averaged Navier-Stokes (RANS) equations, the Navier-Stokes equations for the instantaneous velocity and pressure fields are decomposed into a mean value and a fluctuating component. The averaging process may be thought of as time averaging for steady-state situations and ensemble averaging for repeatable transient situations. The resulting equations for the mean quantities are essentially identical to the original equations, except that an additional term now appears in the momentum transport equation. This additional term is a tensor quantity, known as the Reynolds stress tensor, which has the following definition.

$$
T_{t} \equiv-\rho \overline{v^{\prime} v^{\prime}}=-\rho\left[\begin{array}{lll}
\overline{\frac{u^{\prime} u^{\prime}}{}} & \overline{u^{\prime} v^{\prime}} & \overline{u^{\prime} w^{\prime}}  \tag{3.3}\\
\frac{u^{\prime} v^{\prime}}{u^{\prime} w^{\prime}} & \frac{\overline{v^{\prime} v^{\prime}}}{\overline{v^{\prime} w^{\prime}}} & \frac{\overline{v^{\prime} w^{\prime}}}{w^{\prime} w^{\prime}}
\end{array}\right]
$$

To model the Reynolds stress tensor of above equation in terms of the mean flow quantities，and henceprovide closure of the governing equations，two basic approaches are used in ANSYS FLUENT：

## Eddy viscosity models

Eddy viscosity models use the concept of a turbulent viscosity to model the Reynolds stress tensor asa function of mean flow quantities．The most common model is known as the Boussinesq，approximation：

$$
\begin{equation*}
T_{t_{1}}=2 \mu_{t} s-\frac{2}{3}\left(\mu_{t} \nabla \cdot \vec{v}+\rho K\right) I \tag{3.4}
\end{equation*}
$$

where $\mathbf{S}$ is the strain tensor


And K is the turbulent kinetic energy，$\mu \mathrm{t}$ is the turbulent viscosity．
The eddy viscosity models in ANSYS FLUENT solve additional transport equations for scalar quaditities that enable the turbulent viscosity to be derived．These include Spalart－Allmaras，K－Epsilonand K－Omega models．
＂涫官

## Reypotds Stress Transport Model

Reyriolds stress transport models，also known as second－moment closure models，solve transport equations are solved for each component of the Reynolds stress tensor．

K－Epsilon models provide a good compromise between robustness，computational cost and accuracy．They are generally well suited to industrial－type applications that contain complex recirculation，withor without heat transfer．
$-F$
K－Omega medels are similar to K－Epsilon models in that two transport equations are solved but differin the choice of the second transported turbulence variable．The performance differences are likely tobe a restlt of the subtle differenees in the models，rather than a higher degree of complexity in the physics being captured．These models have seen most application in the aerospace industry．Therefore，they are recommended an alternative to the Spalart－Allmaras models for similar types of applications．

## 3．3．3 METHODOLOGY OF CFD

In general，CFD simulations can be distinguished into three main stages，which are；
1）pre－processor
2）simulator or solver and
3）post－processor

## TRIAL DONE FOR VORTEX SHEDDING

## 1. Pre-Processor

The pre-processing stage includes:

- Solution domain defining. The solution domain defines the abstract environment where the solution is calculated. The shape of the solution domain can be circular or rectangular. Generally, many simulations use a rectangular box shape as the solution domain as shown in the figure.

The choice of solution domain shape and size-can affect the solution of the problem. The smaller sized of domains need less iterations to solve the problem, in contrast to big domains, which need more time to find the solution.



## $\cdots-n^{3}$



Fig. 3.2: Domain condition of the cylinder

| D 1 | 0.90 m |
| :--- | :--- |
| V 2 | 6.3 m |
| H 4 | 6.3 m |
|  | 6.3 m |

Table no 1.: Dimensions of the Domain


## （7）監 合 童



Fig 3．3．：CyTinder with fixed fairing of 0．125D


Fig 3．4：Cylinder fixed with fairing 0．25D

## －Mesh generation

After the solution domain has been defined，we shall generate the mesh within the solution domain．The term mesh generation and grid generation is often interchangeable．The mesh or grid is defined as the discrete locations at which the variables are to be calculated and to be solved．The grid divides the solution domain into a finite number of sub domains，for instance elements，control volumes etc．


Fig 3.5: Meshing for the bare cylinder.


Geometry parts are used only to define the faces, edges, and vertices that make up the surfaces of themodel. The simulation domain on which mesh is generated, and for which physics is solved, is definedusing regions, boundaries, and interfaces. An essential part of the process in setting up your simulationis therefore to define the relationship between geometry parts and regions, boundaries, and interfaces. The mesh continuum models cannot be used for geometry parts. In order to mesh geometry parts, theymust be assigned to one or more regions, and the mesh is then generated for the region.
In this study seven different configuration for the conduits are made and its simulated for 2 different free stream velocities the computational domain for the configuration are explained below: -


Fig 3.8: Domain configuration for cylinder with fairing 0.125 D fixed at different angtes.

Fig 3.9: Dqmain configuration for the cylinder with fairing 0.25D at 180 degree.

## 3.5-BOUNDARY CONDITIONS

There äre several boundary conditions for the discredited equations. Some of them which are used for this work are velocity inlet, pressure outlet, free slip wall, no slip wall.

## Velocity Inlet



A velocity inlet boundary represents the inlet of a duct at which the flow velocity is known... It is applied for inlet to a simulation of incompressible internal flows. It can be used in combination with pressure outlet and flow split outlet boundaries.

## Pressure Outlet

A pressure outlet boundary is a flow outlet boundary at which the pressure is specified. It can apply as an outflow of compressible and internal flows.

## Wall

A wall boundary represents an impermeable surface. Applications of this boundary are as an impenetrable boundary for inviscid flows, as an impenetrable, no-slip boundary for viscous flow, "Slip walls" representing an impenetrable but traction-free surface.

Boundary Conditions Applied.
Inlet - velocity inlet.
Outlet- pressure outlet
Top and bottom - wall (free slip)
cylinder- wall (no slip)

### 3.6 PHYSICS MODELING

This chapter includes the details of physics models selected, solver setup, and stopping criteria. Thischapter also includes details of spring model implemented, Calculatiọn of Eigen frequency of spar still water using free decay test, flow velocities corresponding to the reduced velocities, in lockin rangeis explained.

## Physics models

The physics model-in ANSYS FLUENT defines how a physical phenomenon in a continuay is represented. Essentially, physics models define the primary variables of the simulation (such as pressure, temperature, and velocity) and what mathematical formulation is used to generate the solution. Models in ANSYS FLUENT have varying levels of complexity and functionality but their major purpose are to work with solvers toobtain a solution: Both background region overset region are provided with the same physics continuum. Table gives


The material model is responsible for managing the material, that is, substance or substances, being simulated in the continuum. The material, in turn, is responsible for managing the various thermodynamic and transport properties relevant to that material and to the physical processes being modelled in the continuum. Here in our study the risers are subjected within the sea, so the material chosen here is seawater and the properties are mentioned in the below table.

Table 3 Material Properties of Fluid

| Material | Density | Dynamic Viscosity |
| :--- | :--- | :--- |
| water | $1025 \mathrm{~kg} / \mathrm{m}^{3}$ | $8.887 \mathrm{E}-4$ Pa-s |

### 3.6.2 olution Methods

The Coupled Flow and Coupled Energy models are best for compressible flows, natural convection problems, and flows with large body forces or energy sources. The Segregated Flow model is an alternative for incompressible or mildly compressible flows, particularly when computational resources are an issue; it is used with the Segregated Fluid Energy models. For complex time-periodicflows, you can use Harmonic Balance Flow and Energy. In this study materials water and air are bothprovided considered as incompressible to reduce computational time. So, the segregated flow model is selected which is suitable for incompressible or mildly compressible flows. The Segregated Flow modelsofvés the flow equations (one for each component of velocity, and one for pressure) in a segregated, or uncoupled, manner The linkage between the momentum and continuity equations is achieved with a predictor-eorrector approach. A flow that is in a state of continuous instability, exhibiting irregular, small-scale, high frequency fluctuations in both space and time is termed turbulent.In this simulation since turbulent flow is encountered, turbulent flow viscous regime is selected, And the turbulence model selected is wale sub grid scale large eddy simulation model.

### 3.6.3 Modelling time



The Implicit Unsteady model is the only unsteady model available with the Segregated Flow and -Segregated Fluid Energy models. It uses the Implicit Unsteady solver. Since segregated flow model isselected, implicit unsteady model is selected.

### 3.7 SOLVERS

Solvers control the solution and are activated once per iteration (or once per time-step for explicit unsteady simulations). Typically, a model elects the solvers that are required. Different models can use the same solver and sometimes models need more than one solver. Some models control other models and some only perform task when the solution is initialized. For this reason, not all models elect solvers. Models can also elect solvers that perform subsidiary tasks in addition to the solver thatthe model elected to control it. The following solvers have been used in this study.

### 3.7.1 SōlverParameter



In order to solve the present problem certain parameters need to be fixed so that computational grid can be solved effectively to achieve an approximate solution to the real problem. The solver parameters used in the present investigation is is ted in Table.


| Parameters | Adopted solver settings |
| :--- | :--- |
| Solver | Implicit unsteady |
| Velocity formulation | Absolute |
| Viscous Model | k -omega SST model |
| Operational Pressure | 101325 PA |
| Gravitational Constant | $981 \mathrm{~m} / \mathrm{s}^{2}$ |
| Water Density | Pressure outlet |
| Water Viscosity |  |
| Inlet boundary conditiont |  |
| Cylinder surface bounday_condition | Velocity inlet |
| Outlet boundary condition | Wall (no slip) |

The Implicit Unsteady solver becomes active when the Implicit Unsteady model is active. Its primary fungtion is to control the update at each physical time for the calculation. It also controls the time-stepsize.

## CHAPTER 4

## RESULTS AND DISCUSSION

### 4.1 General

This chapter deals with the discussion about the results of the Numerical simulations performed for bare ylinder and cylifinder with fairings at different angles for different current speed.

### 4.2 CFD Results for the fairings arrangement

### 4.2.1 Bare Cylinder

Case 1). Curtent sped $=0.25 \mathrm{~m} / \mathrm{s}$


Fig 4.1: Bare cylinder velocity Contours for $0.25 \mathrm{~m} / \mathrm{s}$

Fig 4.2: Bare cylinder Coefficient of lift $\mathrm{v} / \mathrm{s}$ flowitimefor $0.25 \mathrm{~m} / \mathrm{s}$
Case 2). Current speed $=0.5 \mathrm{~m} / \mathrm{s}$




Fig 4.4: Bare cylinder Coefficient of lift $\mathrm{v} / \mathrm{s}$ flowtime for $0.5 \mathrm{~m} / \mathrm{s}$

### 4.2.2 Cylinder with fairing of $\mathbf{0 . 1 2 5 D}$ at an angle of Zero degree.

Case 1). Current speed $=0.25 \mathrm{~m} / \mathrm{s}$


Fig 4.5:-Veldcity contours for cyl with fairing 0.125 D at an angle of zero degree for $0.25 \mathrm{~m} / \mathrm{s} \mathrm{CS}$



Fig 4.6: Coefficient of lift $\mathrm{v} / \mathrm{s}$ flowtime for cyl with fairing 0.125 D at an angle of zero degree for $0.25 \mathrm{~m} / \mathrm{s} \mathrm{CS}$

Case 2). Current speed $=0.5 \mathrm{~m} / \mathrm{s}$

Fig 4.7: Velocity contours for cyl with fairing 0.125 D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{s} \mathrm{CS}$



Fig 4.8: Coefficient of lift v/s flowtime for cyl with fairing 0.125 D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{s} \mathrm{CS}$

### 4.2.3 Cylinder with Fairing of 0.125 D fixed at an angle of 45 degree.

Case 1). Current speed $=0.25 \mathrm{~m} / \mathrm{s}$


Fig 4.9:-Velocity Contours for cyl with fairing 0.125 D at an angle of zero degree for $0.25 \mathrm{~m} / \mathrm{s}$ CS





Coser
(1)

Fig 4.10: coefficient of lift $\mathrm{v} / \mathrm{s}$ flowtime for cyl with fairing 0.125 D at an angle of zero degree for $0.25 \mathrm{~m} / \mathrm{s} \mathrm{CS}$
Case 2). Current speed $=0.5 \mathrm{~m} / \mathrm{s}$

Fig 4.11: Velocity dontours for cyl with fairing 0.125 D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{s} \mathrm{CS}$
lift-rplot (Time=2.5000e+02 s)


$$
- \text { lift }
$$

Fig 4.12: Coefficient of lift v/s flowtime cyl with fairing 0.125 D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{s} \mathrm{CS}$

### 4.2.4 Cylinder with Fairing of 0.125 D fixed at an angle of $\mathbf{9 0}$ degree.

Case 1). Current speed $=0.25 \mathrm{~m} / \mathrm{s}$


Fig 4, +3 : Velocity contours for cyl with fairing 0.125 D at an angle of zero degree for $0.25 \mathrm{~m} / \mathrm{s}$ CS

lift-rplot


Fig 4.14: Coefficient of lift v/s flowtime for cyl with fairing 0.125 D at an angle of zero degree for $0.25 \mathrm{~m} / \mathrm{s}$ CS

Case 2). Current speed $=0.5 \mathrm{~m} / \mathrm{s}$


Fig 4.15: Velocity dontours for cyl with fairing 0.125 D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{s} \mathrm{CS}$


Fig 4.16: Coefficient of lift vfs flowtime for cyl with fairing 0.125D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{sCS}$

### 4.2.5 The cylinder with fairing of 0.125 D fixed at an angle of $\mathbf{1 3 5}$ degree.

Case 1). Current speed $=0.25 \mathrm{~m} / \mathrm{s}$


Fig 4.18: Coefficient of lift $\mathrm{v} / \mathrm{s}$ flowtime for cyl with fairing 0.125 D at an angle of zero degree for $0.25 \mathrm{~m} / \mathrm{s}$ CS
Case 2). Current speed $=0.5 \mathrm{~m} / \mathrm{s}$


Fig 4.19. Velocity contours for cyl with fairing 0.125 D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{s}$ CS



Fig 4.20: Coefficient of lift v/s flowtime for cyl with fairing 0.125D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{s} \mathrm{CS}$

### 4.2.6 The cylinder with fairing of $\mathbf{0 . 1 2 5 D}$ fixed at an angle of 180 degree.

Case 1). Current speed $=0.25 \mathrm{~m} / \mathrm{s}$


Fig 4.21:Velocity Contours for cyl with fairing 0.125D at an angle of zero degree for $0.25 \mathrm{~m} / \mathrm{s}$ CS ar
lift-rplot (Time=2.5000e+02 s)

$\square$

- lift

Fig 4.22: Coefficient of lift $\mathrm{v} / \mathrm{s}$ flowtime for cyl with fairing 0.125 D at an angle of zero degree for $0.25 \mathrm{~m} / \mathrm{s} \mathrm{CS}$
Case 2). Current speed $=0.5 \mathrm{~m} / \mathrm{s}$


Fig 4.23. Velocity contours for cyl with fairing 0.125 D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{s} \mathrm{CS}$


Fig 4.24: Coefficient of lift $\mathrm{v} / \mathrm{s}$ flowtime for cyl with fairing 0.125 D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{s}$ CS

### 4.2.7 The Cylinder with fairing of 0.25 D fixed at an angle of $\mathbf{1 8 0}$ degree.

Case 1). Current speed $=0.25 \mathrm{~m} / \mathrm{s}$


Fig 4.26: Coefficient of lift $\mathrm{v} / \mathrm{s}$ flowtime cyl with fairing 0.25 D at an angle of zero degree for $0.25 \mathrm{~m} / \mathrm{s}$ CS
Case 2). Current speed $=0.5 \mathrm{~m} / \mathrm{s}$


Fig 4.27: Velocity contours for cyl with fairing 0.25 D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{s}$ CS


Fig 4.28: Coefficient of lift v/s flowtime cyl with fairing 0.25 D at an angle of zero degree for $0.5 \mathrm{~m} / \mathrm{s}$ CS

## RESULT TABLE

Table No. 5 List of obtained coefficient of lift values.

| Sl. No. | Type of cylinder | Angle | Current speed (m/s) | Lift Coefficient (Cl) |
| :---: | :---: | :---: | :---: | :---: |
| 01 | Bare | - | 0.25 | 0.1913959 |
| 02 | Bare . |  | 005 $0_{-7}$ | 0.5372883 |
| 03 | With 0.125 D fairifíg | Zero | $\overline{0}, 25$ | 0.0675923 |
| 04 | With oi25D fairing | Zero | 0.5 | $1.450176 \mathrm{e}-27 \approx 0$ |
| $05^{\text {a }}$ | With 0.125D fairing | $45^{\circ}$ | 0.25 | $-1.866447$ |
| 06 -7 | With 0.125D fairing | $45^{\circ}$ | 0.5 | $-0.0684217$ |
| $07$ | With 0.125D fairing | $90^{\circ}$ | 0.25 | 0.1640284 |
| $08$ | With 0.125D fairing | $90^{\circ}$ | 0.5 | -5.723619 |
| $0^{3} 9$ | With 0.125D fairing | $135^{\circ}$ | 0.25 | -0.7536188 |
| 10 | With 0.125D fairing | $135^{\circ}$ | 0.5 | -3.105413 |
| 11 | With 0.125D fairing | $180^{\circ}$ | 0.25 | -0.05476332 |
| 12 | With 0.125D fairing | $180^{\circ}$ | 0.5 | 0.1240201 |
| $13$ | With 0.25D fairing | $180^{\circ}$ | 0.25 | -0.01796491 |
| $14$ | With 0.25D fairing | $180^{\circ}$ | 0.5 | -0.6813356 |

## CHAPTER 5

## SUMMARY AND CONCLUSIONS

### 5.1 Summary

In this thesis, numerical simulations of flow past circular conduits were carried out at different free stream velocities for different arrangement of cylinder with fairings fixed at different angles. ANSYS FLUENT was used for this study as a solver for the study. Simulation was carried out over a flow past a circular conduit for two different current speeds. In these different configurations of the conduits the effect of the presence of one another was analysed such as lift variation over time over the domain.

### 5.2 Conclusions

The salient conclusions from the study are as follows:

1. It has been observed that for most cases the values of lift force is higher in case of the bare cylinders in comparison with the results for cylinder arrangement with attached fairing.
2. It has been observed that the lift force is increased for bare cylinder as we increase the current speed.
3. It has been observed that the vortex induced vibrations are reduced and lift force is decreased after attaching the fairings to the cylinder.
4. The trial has been done for two types of fairings i.e, fairing of 0.125 D and fairing of 0.25 D .
5. Thus, by the analysis conducted we can conclude that foricylinder with fairing 0.125D fixed at an angle of $0^{\circ}$ for the current speed $0.5 \mathrm{~m} / \mathrm{s}$ the coefficient of lift is more or less its equal to Zero (0).

## REFERENCES



1. Xiaojian Jin, Shunqi Yang, Hai Chen (2017) Research on Fairing design and CFD Analysis-of Submarine Pipeline-Inspection ARV, MATEC Web of Conferences 108, 07003.
2. Syamsul/Azry Md Esa, Wan Mohd. Arif Aziz Japar, Nor Azwadi Che Sidik (2019) Design and Arralysis of Vortex Induced Vibration (VIV) Suppression Device, Akademia baru ISSN: 2180-1363.
3. Narsahliza Binti Muhamat Yain (2017) Marine riser Vortex Induced Vibration (VIV) suppression Bevice.
$4 . \overline{\text { Syazwan }} \overline{\text { Bin }}$ Sazali (2016) Vortex induced vibration (VIV) reduction measure on dimpled concrete coated pipeline.
4. Satya Prakash Singh, Gautam Biswas, Perumal Nithiarasu (2013) A numerical study of vortex shedding from a circular cylinder vibrating in the in-line direction, DOI: 10.1108/HFF-08-2012-0183.
5. Meghanadhan C A, Sunil A S (2016) An Experimental Study of Vortex Shedding Behind a Bluff Bedy in a Water Charnnel, (IJERT) ISSN: 2278-0181.
6. Alok Mishra, Ashoke De, Suppression of vortex shedding using a slit through the circular cylinder at low Reynolds number.
7. Aswin Janardhanan (2014) Reducing Vortex-Induced Vibration of Drilling Risers yith-Matrine Fairing, Spring 5-16-2014,

